



CHARA Classic/Climb Numbers.





















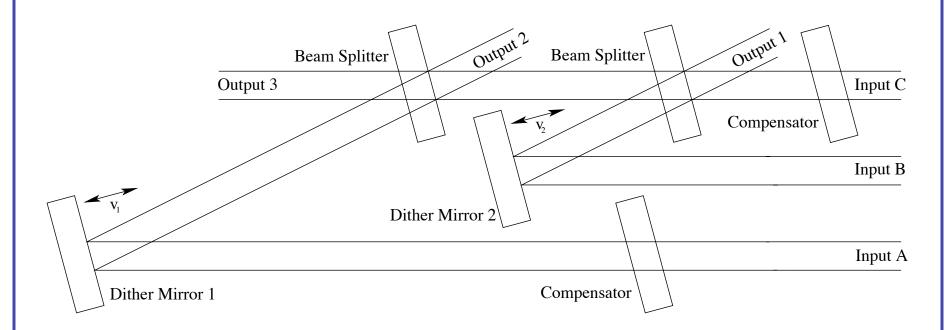






























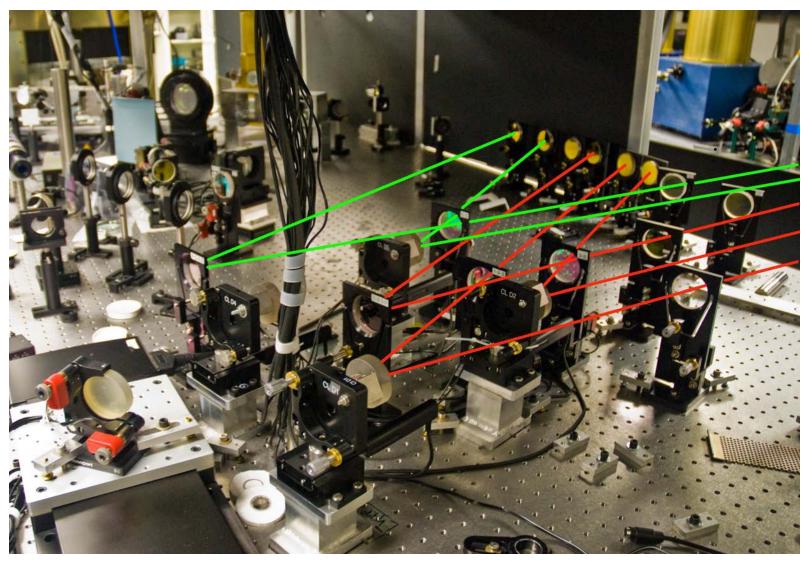








































Automated Data Reduction

- Automated editing Fringe > 1.1 Noise Power
- Took approximately 200 minutes to crunch.
- V < 0 and V > 1 thrown away.
- Not reliable for science.
- K&H magnitudes extracted from 2MASS.
- Stars without 2MASS data thrown away.
- Includes both calibrators and science targets.

















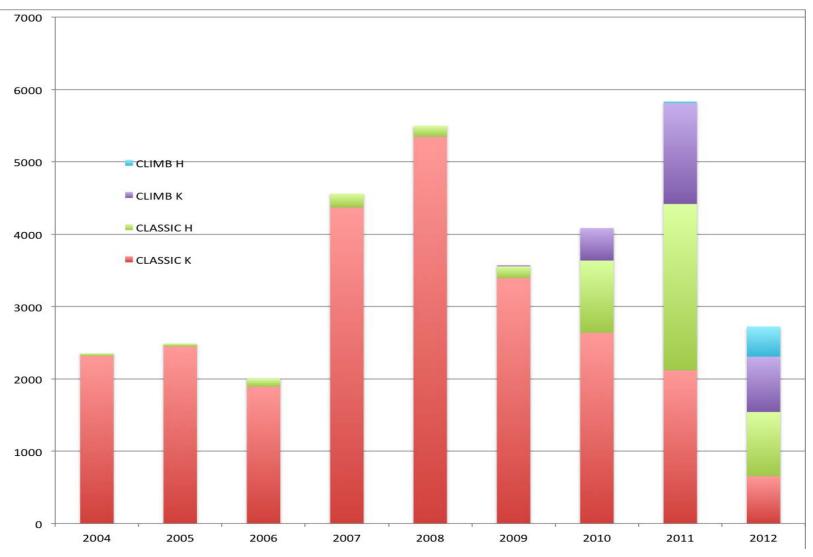








Amount of Data





























K/H Mags are converted to a photon count.

- Uses numbers from Camping, Rieke & Lebofsky PASP 90, 896i (1995): For Mag 0 Star:
 - J Band 1.26 micron: 1603 Jy
 - H Band 1.60 micron: 1075 Jy
 - K Band 2.22 micron: 667 Jy
- 1 Jy = 1.51 x 10⁷ Photons S⁻¹ m⁻² $(d\lambda/\lambda)^{-1}$
- All data are calibrated to 1 second.
- This assumes the NIRO readout mode behaves.
- Camera Gain = 0.3, DQE = 60%.

















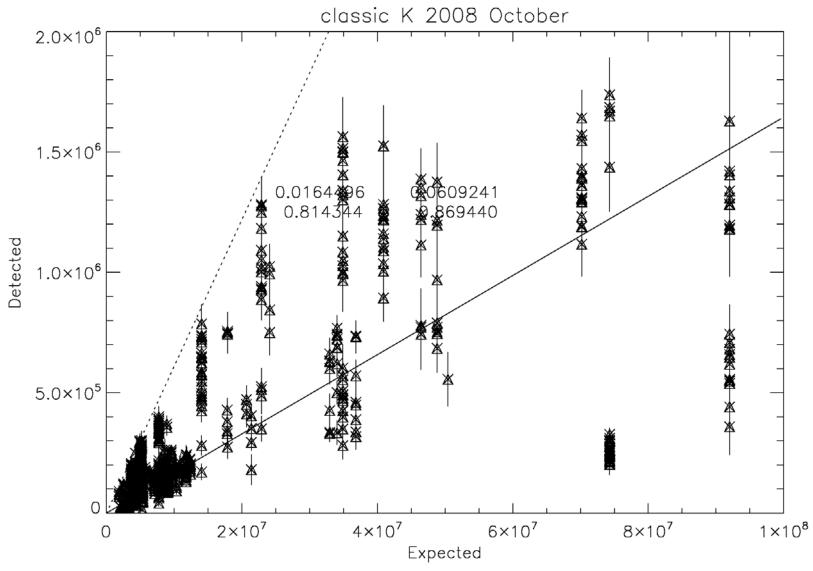
































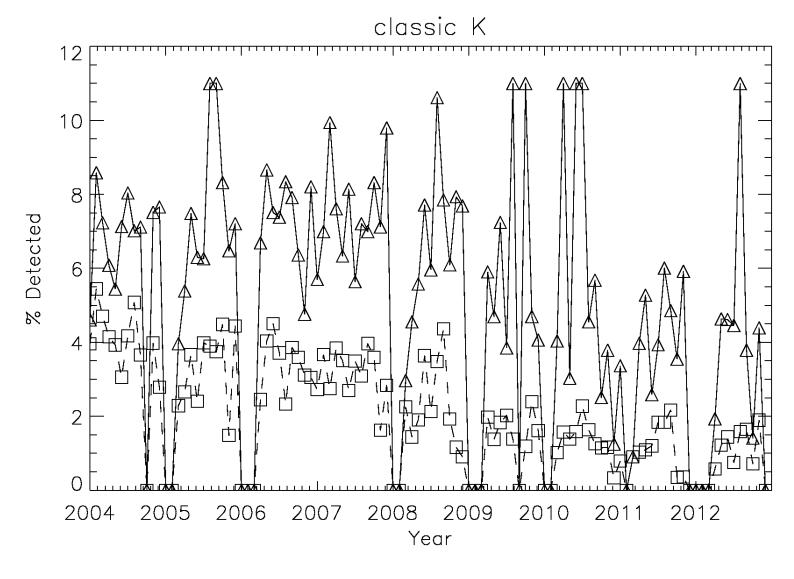
































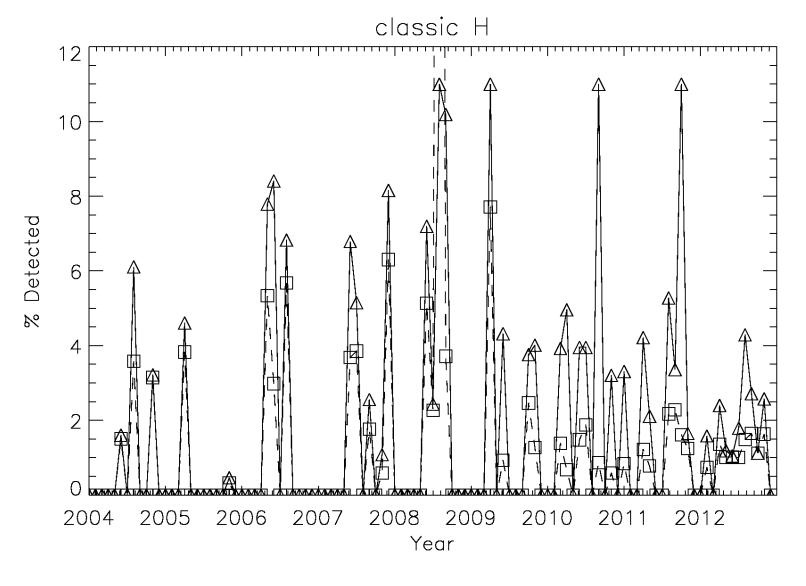
































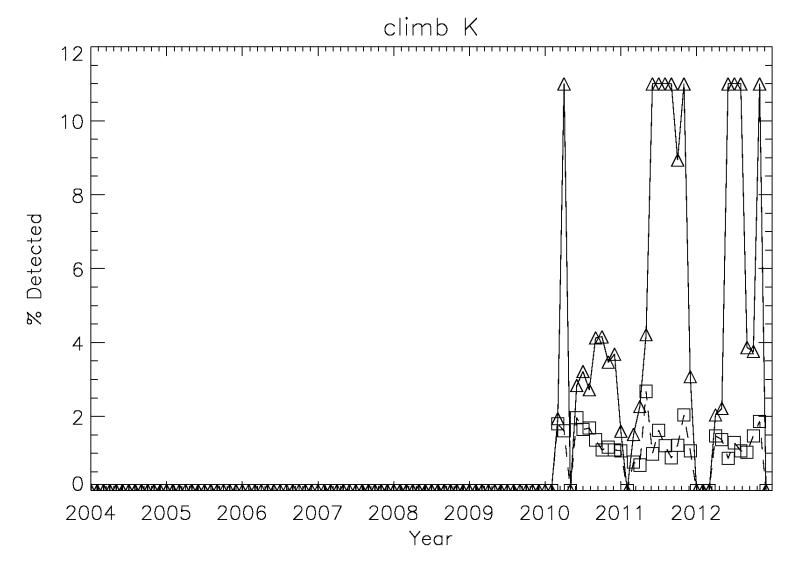
































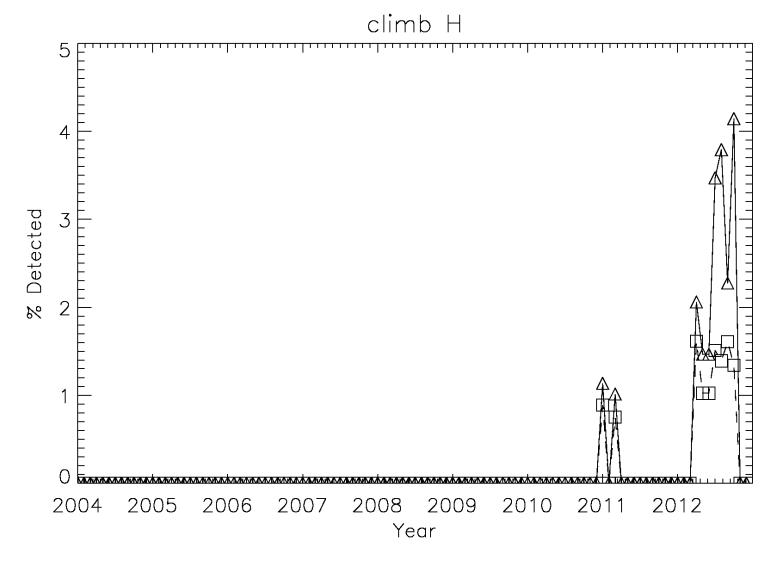






































System Visibilities are Corrected For Estimated Stellar Diameter

A&A 426, 297–307 (2004) DOI: 10.1051/0004-6361:20035930 © ESO 2005 Astronomy Astrophysics

The angular sizes of dwarf stars and subgiants

Surface brightness relations calibrated by interferometry*

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Abstract. The availability of a number of new interferometric measurements of Main Sequence and subgiant stars makes it possible to calibrate the surface brightness relations of these stars using exclusively direct angular diameter measurements. These empirical laws make it possible to predict the limb darkened angular diameters θ_{LD} of dwarfs and subgiants using their dereddened Johnson magnitudes, or their effective temperature. The smallest intrinsic dispersions of $\sigma \le 1\%$ in θ_{LD} are obtained for the relations based on the K and L magnitudes, for instance $\log \theta_{LD} = 0.0502 (B - L) + 0.5133 - 0.2 L$ or $\log \theta_{LD} = 0.0755 (V - K) + 0.5170 - 0.2$ K. Our calibrations are valid between the spectral types A0 and M2 for dwarf stars (with a possible extension to later types when using the effective temperature), and between A0 and K0 for subgiants. Such relations are particularly useful for estimating the angular sizes of calibrators for long-baseline interferometry from readily available broadband photometry.























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Observatoire de la Côte d'Azur, BP 4229, 06304 Nice Cedex 4, France

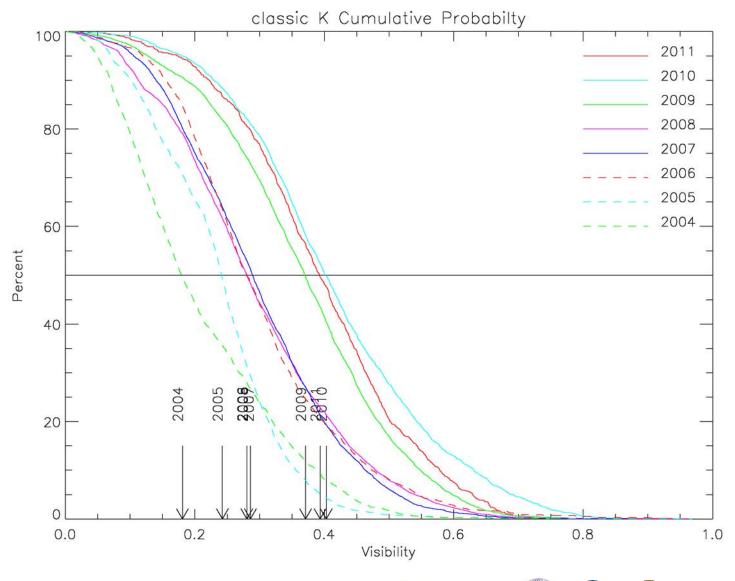
⁴ European Southern Observatory, Karl-Schwarzschild-str. 2, 85748 Garching, Germany

Observatoire de Genève, 1290 Sauverny, Switzerland





Uncorrected result from last year:





















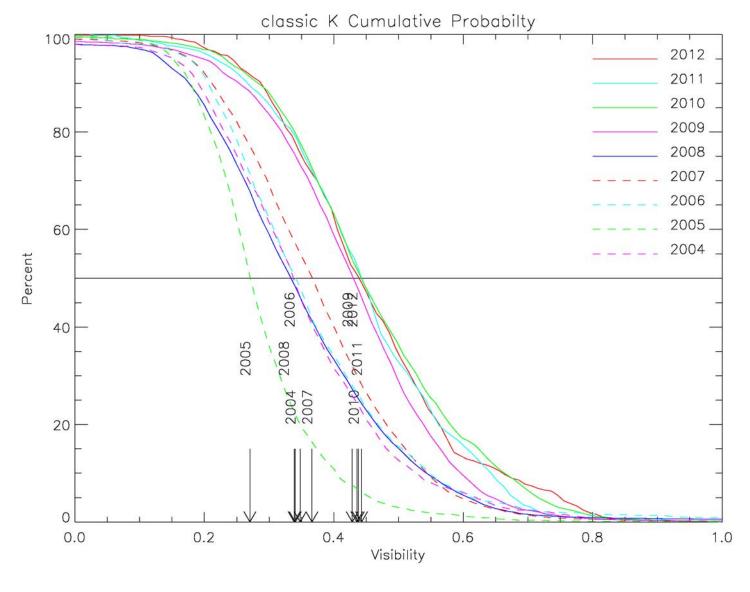
































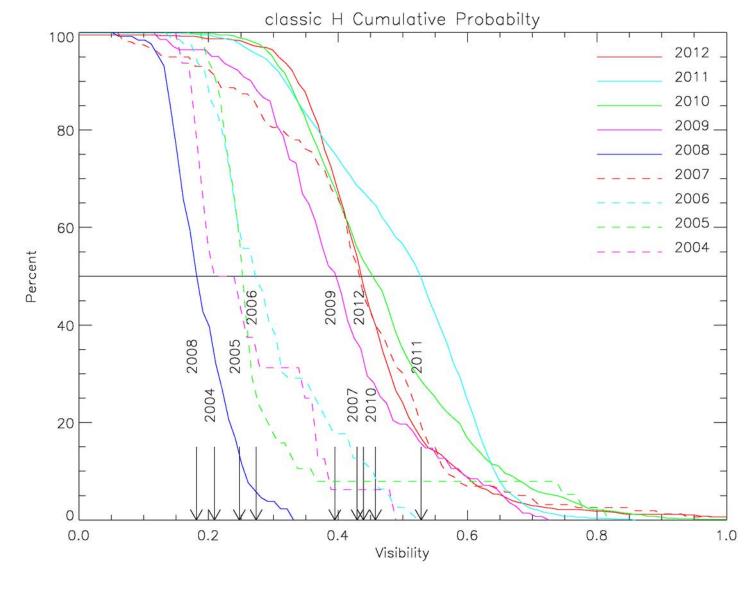






























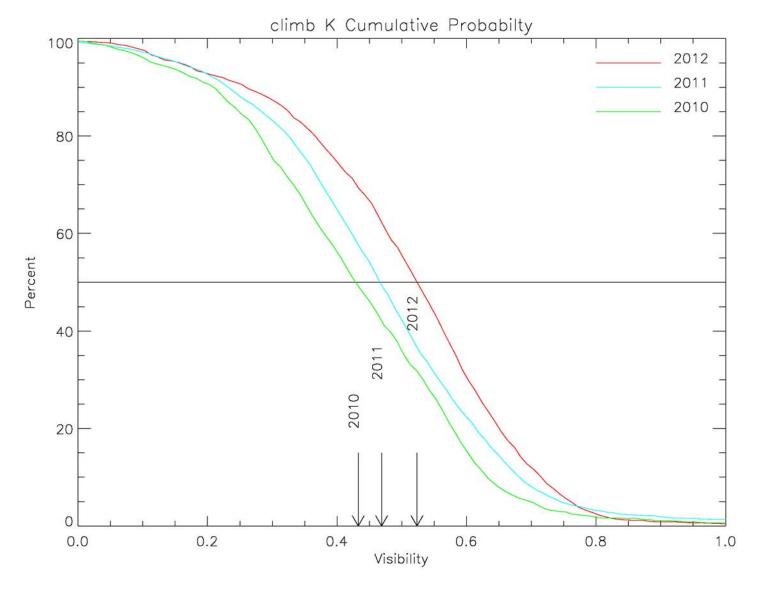
































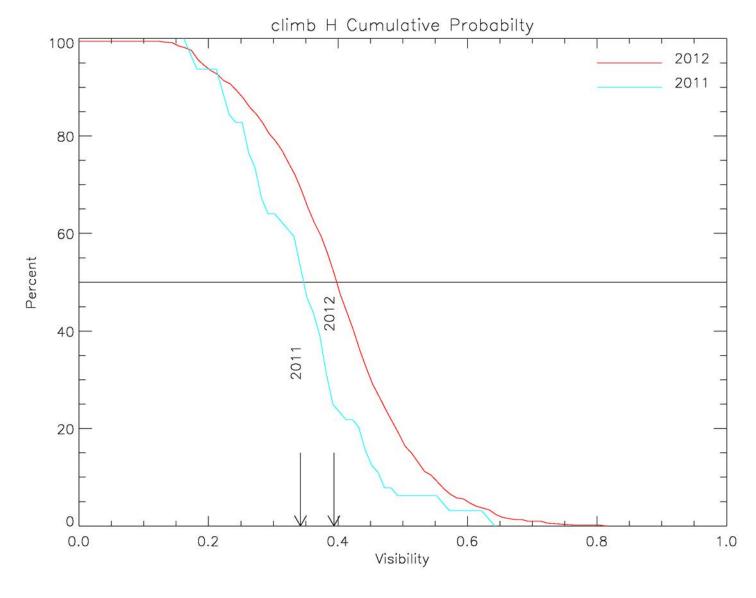




























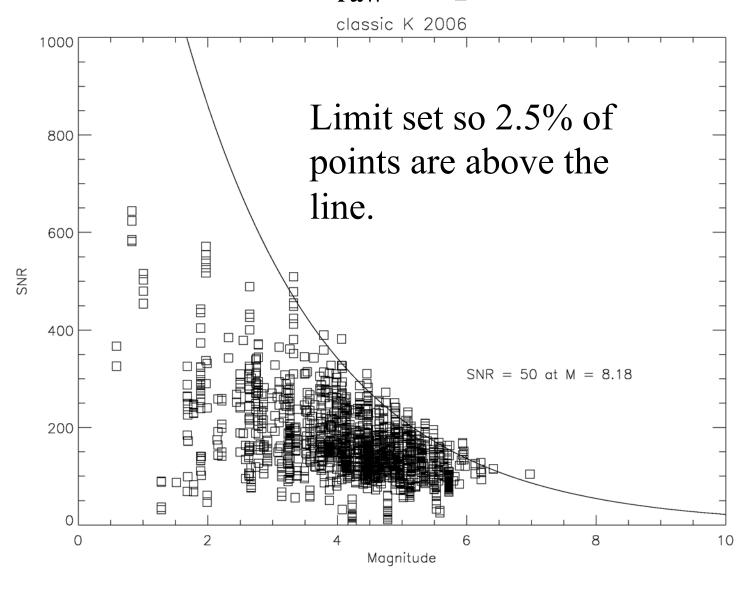








$SNR \sim V_{raw} * sqrt(N)$



















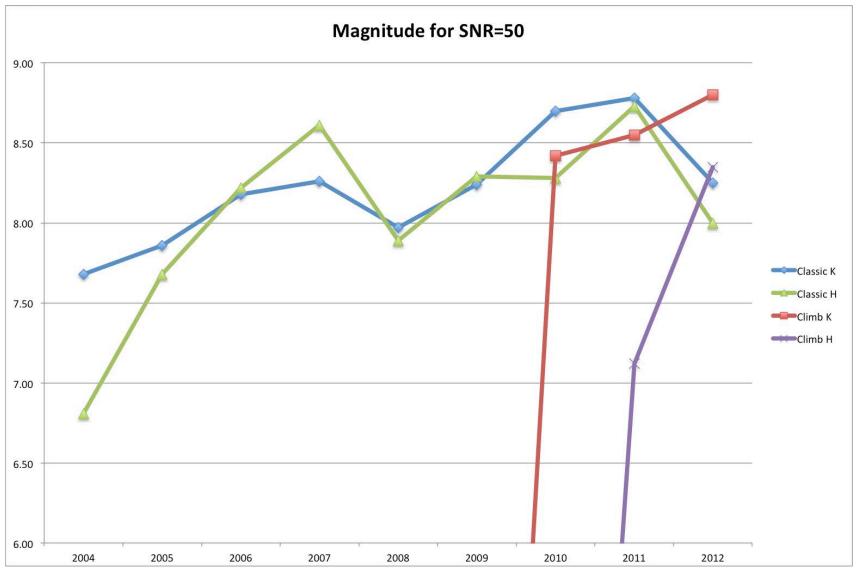






































New Baseline Solution



- The system records the OPLE demand positions and current scope Alt/Az when fringes are found.
- This is now automated for all beam combiners.
- 134889 baseline solution data points were recorded.
- The demand position is better for modeling than the measured position.
- The height of a scope is degenerate with its internal path.
- We use a different internal path for each POP configuration to solve for telescope positions.
- We then do a separate solution for internal path.

















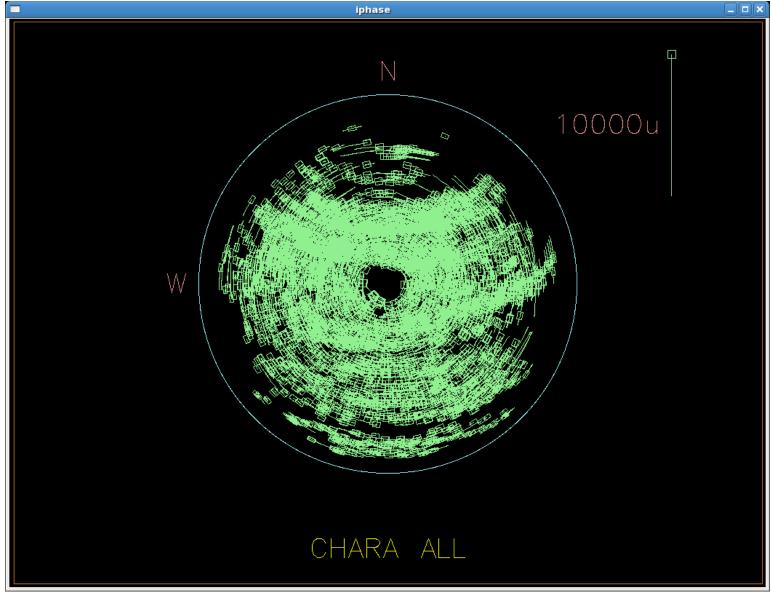






























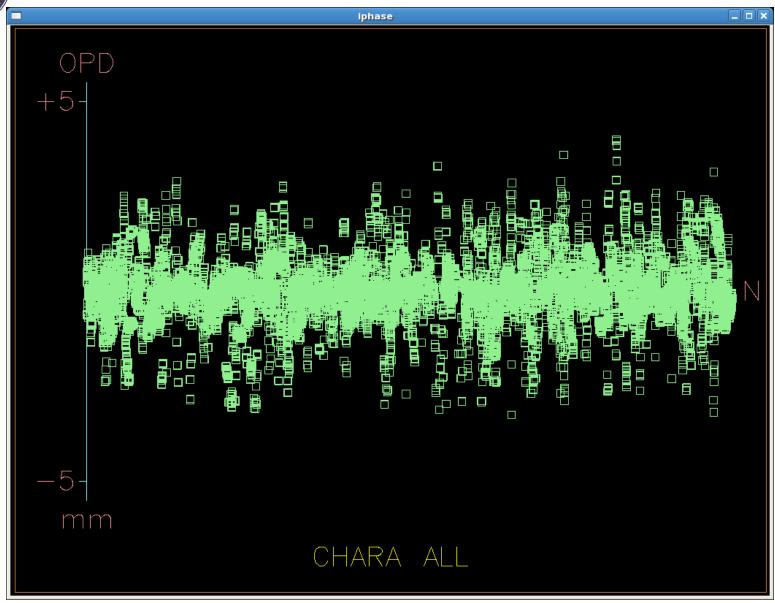




































For telescope S1: value stddev delta (total delta 0.000)

X0FFSET 0.000 0.000 0.000

0.000 0.000 0.000 Y0FFSET

0.000 0.000 0.000 Z0FFSET

0.000 0.000 0.000 LIGHT

For telescope S2: value stddev delta (total delta 1674.412)

X0FFSET -5746856.890 62.733 235.383

YOFFSET 33581561.446 109.897 -1408.412

ZOFFSET 635519.155 163.617 -874.429

LIGHT 4092893.853 2837.526 9328.161

For telescope E1: value stddev delta (total delta 3444.169)

XOFFSET 125332477.873 45.552 1215.095

Y0FFSET 305934252.825 78.549 -1985.452

ZOFFSET -5911377.995 112.808 2538.469

LIGHT 11259794.703 1609.680 -3527.750

For telescope E2: value stddev delta (total delta 1887.612)

XOFFSET 70394847.986 67.731 1182.487

YOFFSET 269715332.908 95.540 -1446.305

ZOFFSET -2799046.780 144.896 270.192

LIGHT 22697827.994 1782.246 -2670.726

For telescope W1: value stddev delta (total delta 2121.951)

XOFFSET -175071927.982 55.365 -478.712

YOFFSET 216320939.168 76.681 -1994.939

ZOFFSET -10791048.377 123.469 541.966

LIGHT 27290684.165 1918.332 -4125.126

For telescope W2: value stddev delta (total delta 2603.854)

XOFFSET -69091262.937 58.622 -1565.808

YOFFSET 199335418.222 84.613 -1571.223

465199.124 153.986 1363.657 Z0FFSET

LIGHT -10864114.917 1829.877 -5745.406



















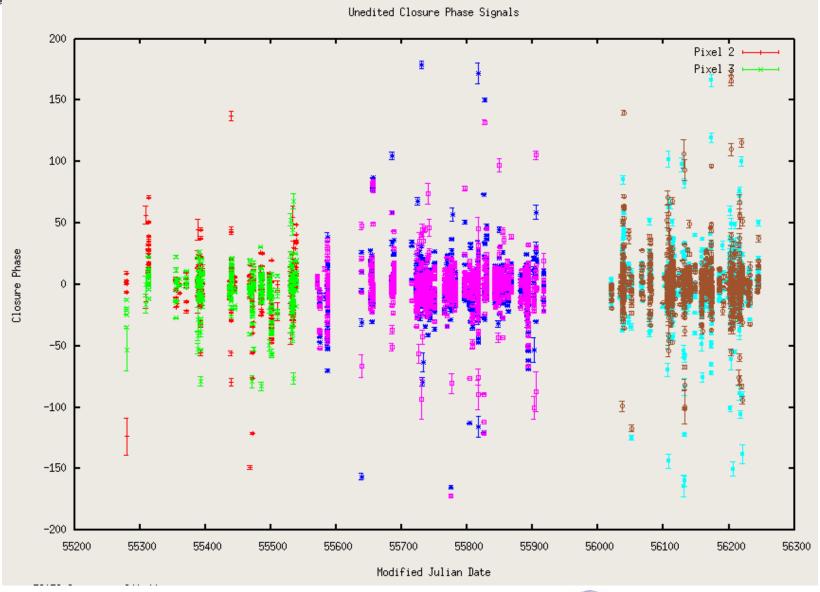




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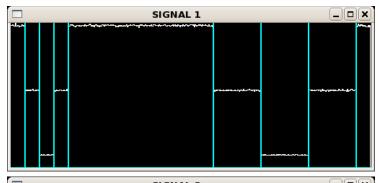


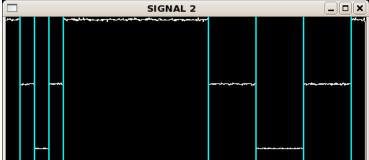


Old Shutter Sequence

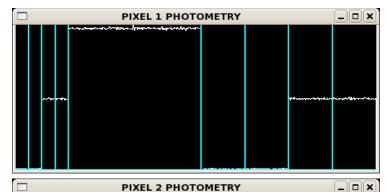


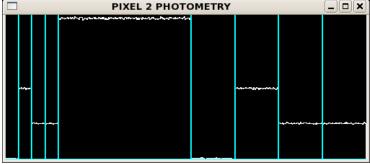
Classic

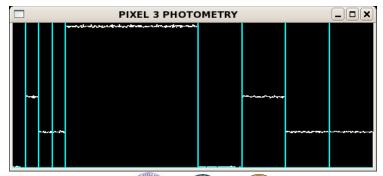




Climb

































This made the Math difficult:

Classic

For each part of the shutter sequence, beam A, no beams and beam B, we do exactly the same analysis as described above for the fringe data. This results in the three power spectra

$$PS(f_A(t)) = PS(\Im_A(t)) + PS(P_A(t)) + PS(C(t))$$
(59)

$$PS(f_B(t)) = PS(\Im_B(t)) + PS(P_B(t)) + PS(C(t))$$
(60)

$$PS(f_{Dark}(t)) = PS(C(t))$$
(61)

where here $\Im_i(t)$ is the scintillation in the signal with just input beam i, $P_i(t)$ is the photon noise with just input beam i, and C(t) is the camera noise. We expect the noise in the output signal to contain the scintillation and photon noise from both input beams, but to have the camera noise just once. We therefore write the noise power spectrum in the signal to be:

$$PS(\aleph(t)) = PS(f_A(t)) + PS(f_B(t)) - PS(f_{Dark}(t))$$
(62)

and this must be subtracted from the fringe signal power spectrum of Equation 59

$$PS(f(t)) = PS(f_{norm}(t)) - PS(\aleph(t))$$
(63)

and we can now perform the analysis set out in section 2.4.2.





























This made the Math difficult:

Climb

It is now possible to calculate the noise power in the fringe power spectra of Equations 66, 67, and 68:

$$PS(\aleph_{AB}(t)) = \left(PS(\aleph_2(t))/T_{2AB}^2 + PS(\aleph_3(t))/T_{3AB}^2\right)/4 \tag{76}$$

$$PS(\aleph_{BC}(t)) = \left(PS(\aleph_1(t))/T_{2BC}^2 + (PS(\aleph_2(t)) + PS(\aleph_3(t)))/(T_{2BC}^2 + T_{3BC}^2)\right)/4 (77)$$

$$PS(\aleph_{CA}(t)) = \left(PS(\aleph_2(t))/T_{2CA}^2 + PS(\aleph_3(t))/T_{3CA}^2\right)/4.$$
(78)

These must be subtracted from the fringe signals to obtain the noise free signals

$$PS(f_{AB}(t)) = PS(f_{AB,\text{norm}}(t)) - PS(\aleph_{AB}(t))$$
(79)

$$PS(f_{BC}(t)) = PS(f_{BC,norm}(t)) - PS(\aleph_{BC}(t))$$
(80)

$$PS(f_{CA}(t)) = PS(f_{CA,\text{norm}}(t)) - PS(\aleph_{CA}(t))$$
(81)

suitable for the analysis set out in section 2.4.2.























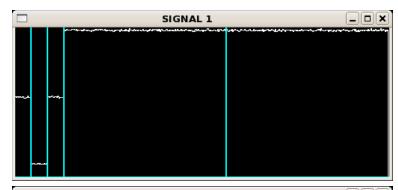




New Shutter Sequence



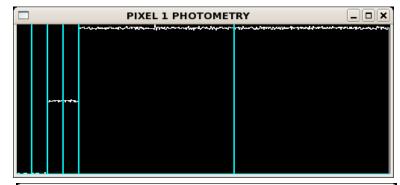
Classic

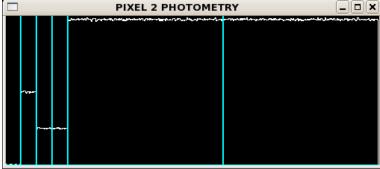


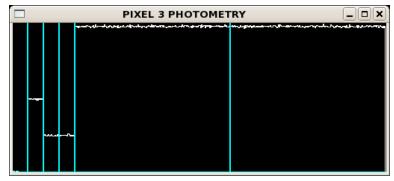


This gives a direct measurement of the background power.

Climb































CHARA Rocks!

























